

Through the looking glass: And what nucleons do there

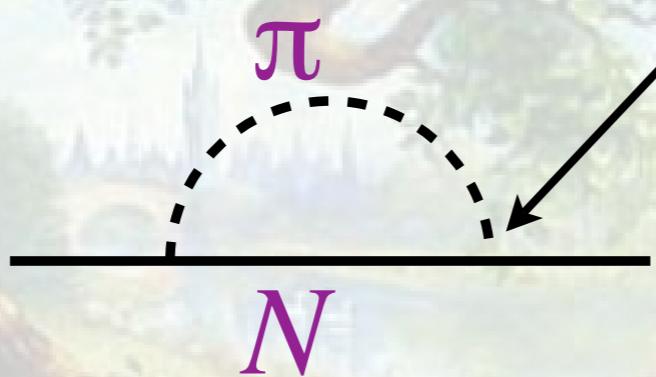
*HUGS, 5-23 June 2006
Lecture 2/2*

*Ross Young
Jefferson Lab*

Chiral correction

$$O(p^2) : \quad M_0 + c_2 m_\pi^2$$

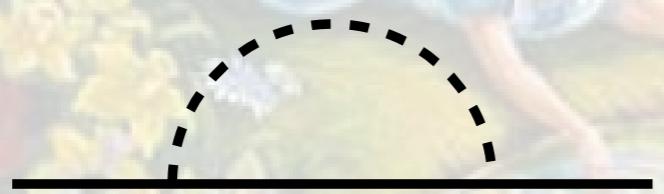
$$O(p^3) :$$



$$\sim \frac{g_A}{f_\pi} \quad \begin{array}{l} \text{Nucleon axial charge} \\ \text{Pion decay constant} \end{array}$$

*Quantum fluctuation, integrate over
all intermediate pion energies*

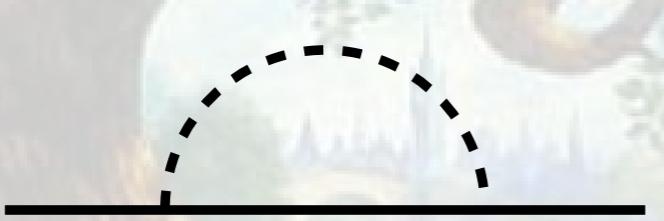
*Do some spin algebra, isospin sum, heavy-nucleon limit,
and temporal and angular integration:*



$$= -\frac{3g_A^2}{16\pi^2 f_\pi^2} \int_0^\infty dk \frac{k^4}{k^2 + m_\pi^2}$$

Poor-man renormalisation

$$O(p^2) : \quad M_0 + c_2 m_\pi^2$$


$$= -\frac{3g_A^2}{16\pi^2 f_\pi^2} \int_0^\infty dk \frac{k^4}{k^2 + m_\pi^2}$$

cubic divergence!

$$\begin{aligned} \int_0^\infty dk \frac{k^4}{k^2 + m_\pi^2} &= \int_0^\infty dk \frac{k^4 - m_\pi^4 + m_\pi^4}{k^2 + m_\pi^2} \\ &= \int dk (k^2 - m_\pi^2) + \textcircled{ } \int dk \frac{m_\pi^4}{k^2 + m_\pi^2} = \textcircled{ } \frac{\pi}{2} m_\pi^3 \end{aligned}$$

$$M_N = M_0 + c_2 m_\pi^2 - \frac{3g_A^2}{16\pi^2 f_\pi^2} \left[\int dk k^2 - m_\pi^2 \int dk + \frac{\pi}{2} m_\pi^3 \right]$$

Renormalised expansion

$$M_N = M_0 + c_2 m_\pi^2 - \frac{3g_A^2}{16\pi^2 f_\pi^2} \left[\int dk k^2 - m_\pi^2 \int dk + \frac{\pi}{2} m_\pi^3 \right]$$

$$M_0^{ren} = M_0 - \frac{3g_A^2}{16\pi^2 f_\pi^2} \int dk k^2$$

$$c_2^{ren} = c_2 + \frac{3g_A^2}{16\pi^2 f_\pi^2} \int dk$$

*Absorb infinities into
redefinition of expansion
constants*

$$M_N = M_0^{ren} + c_2^{ren} m_\pi^2 - \frac{3g_A^2}{32\pi f_\pi^2} m_\pi^3$$

Nonanalytic term: model-independent

Chiral structure of the nucleon

*Chiral corrections to electromagnetic
form factors*

Leading nonanalytic corrections

$$\mu^{p(n)} \sim \mp \frac{g_A^2 M_N}{8\pi f_\pi^2} m_\pi$$

$$\langle r^2 \rangle_E^{p(n)} \sim \mp \frac{1 + 5g_A^2}{(4\pi f_\pi)^2} \log \frac{m_\pi}{\Lambda_\chi}$$

*charge radii become
infinite in the chiral limit*



The Foldy Term

$$G_E(Q^2) = F_1(Q^2) - \frac{Q^2}{4M^2} F_2(Q^2)$$

$$G_M(Q^2) = F_1(Q^2) + F_2(Q^2)$$

$$F_1^n(Q^2 = 0) = 0 \quad \mu^n \equiv G_M^n(Q^2 = 0) = F_2^n(0)$$

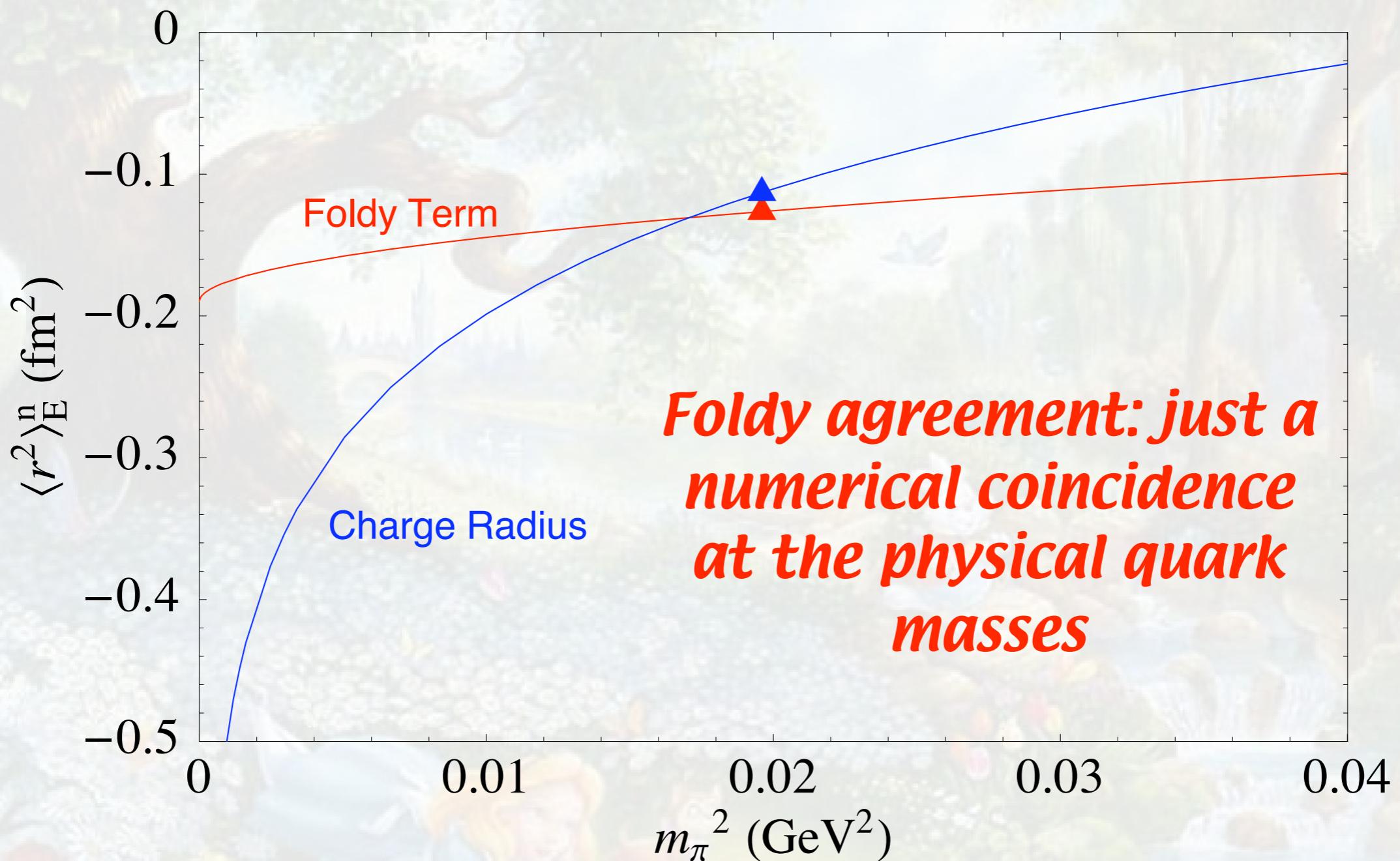
$$\langle r^2 \rangle_E^n = -6 \frac{d}{dQ^2} G_E^n(Q^2 = 0) = -6 \frac{d}{dQ^2} F_1^n(0) + \frac{3}{2M^2} \mu^n$$

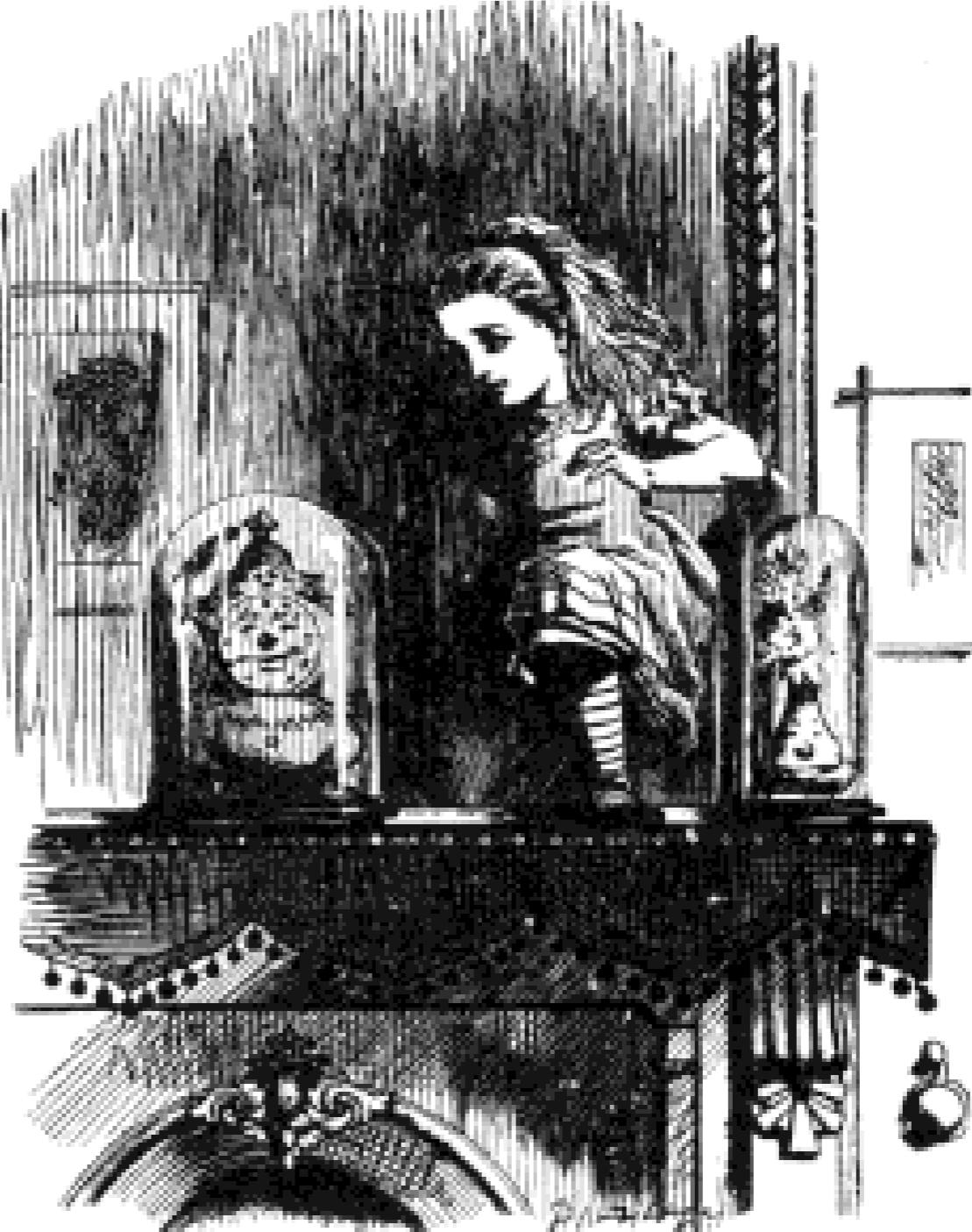
$$\langle r^2 \rangle_E^n = -0.116 \pm 0.002 \text{ fm}^2$$

-0.126 fm²

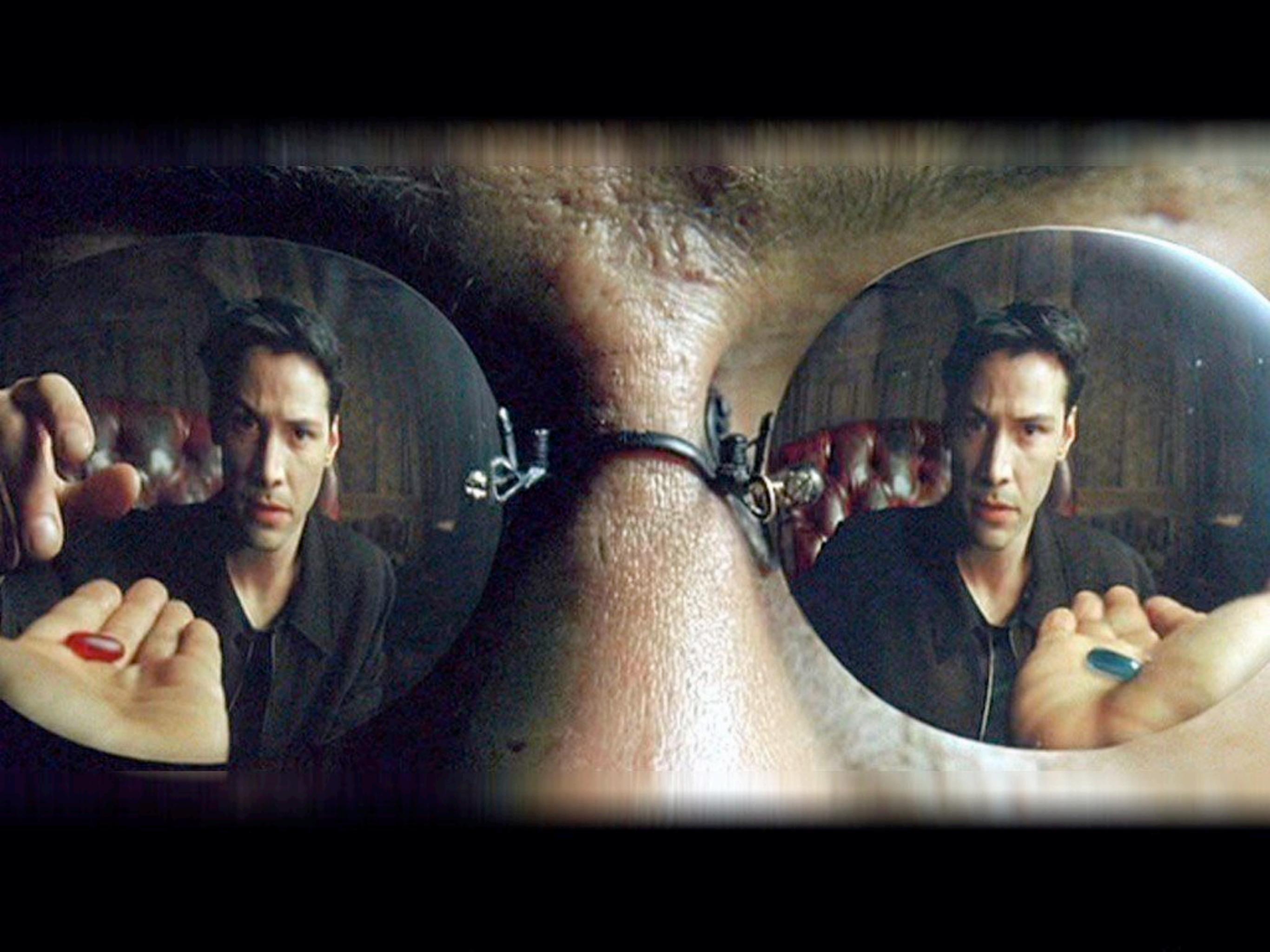
Intrinsic charge distribution interpreted to be zero, charge radius just a relativistic effect of anomalous magnetic moment

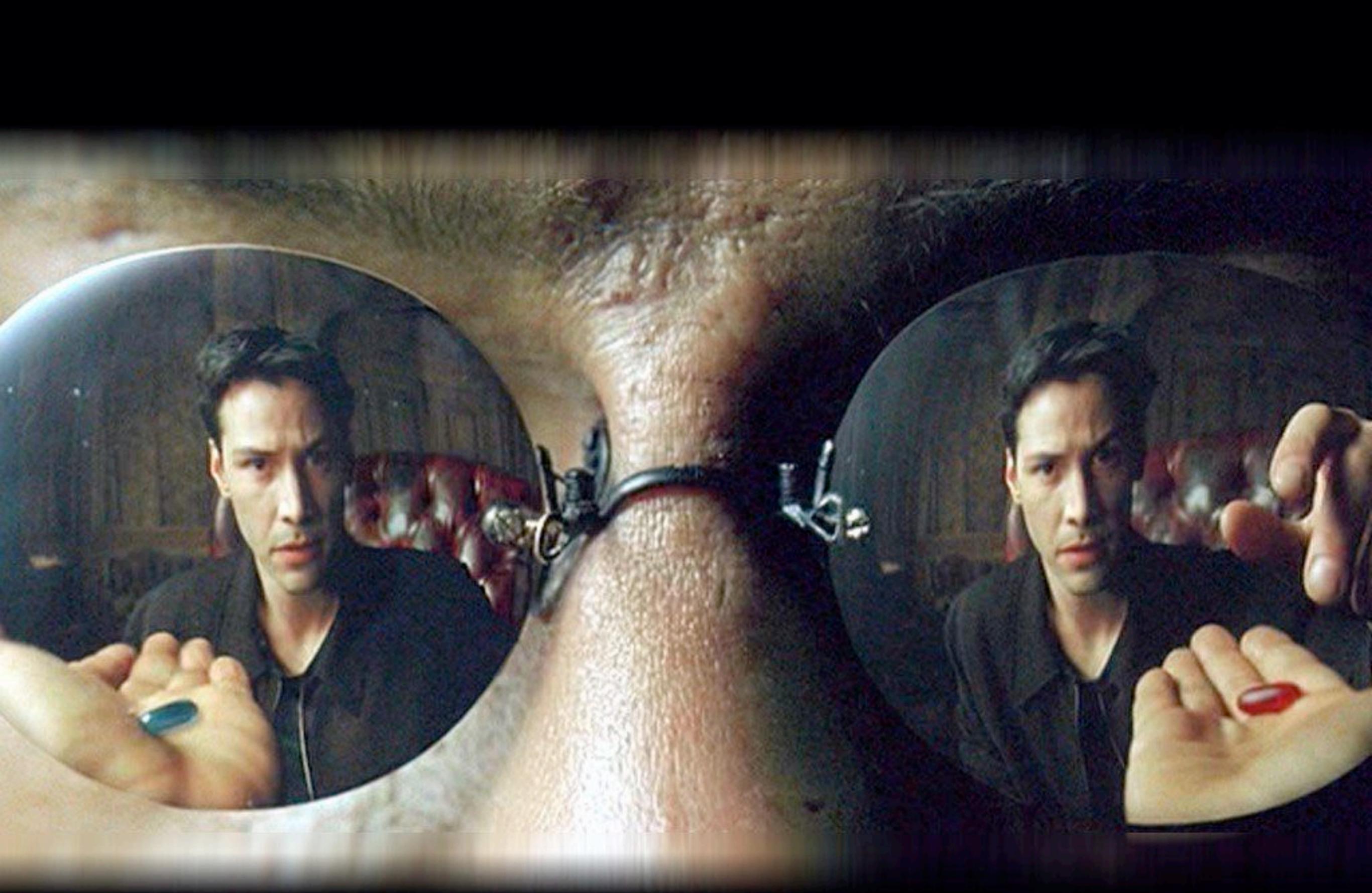
Foldy Term (1952)





Entering the parity-violating world

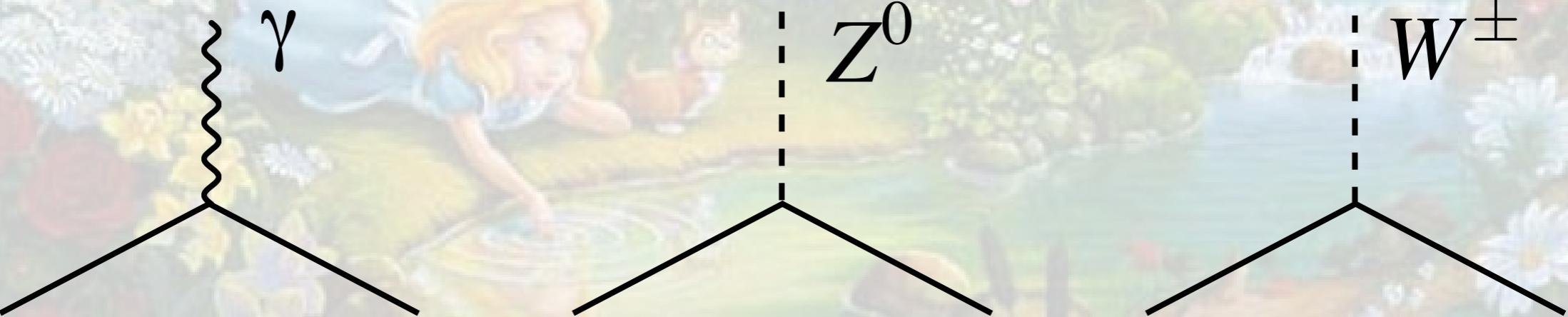




The weak-interaction

- Standard Model electroweak theory unified by Weinberg, Salam & Glashow
 - Local gauge symmetry group $SU(2) \times U(1)$
 - Parity maximally broken by $SU(2)$ sector, coupling to left-handed currents only

$$\mathcal{L}_{\text{int}} = -e J_{EM}^\mu(x) \mathcal{A}_\mu(x) - \frac{g}{2 \cos \theta_W} J_{NC}^\mu(x) Z_\mu(x) - \frac{g}{2\sqrt{2}} J_{CC}^\mu(x) W_\mu^\dagger(x) + \text{h.c.}$$

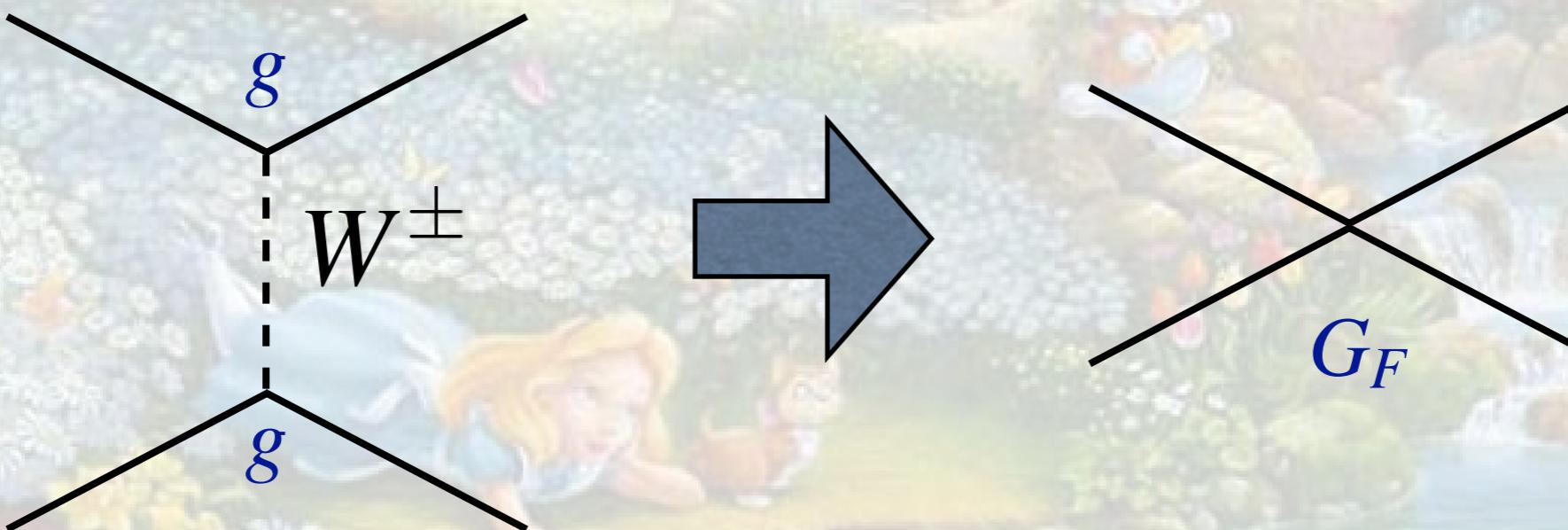


The weak-interaction... 2

- *Weak mixing angle*
 - *Physical photon is mixture of U(1) and SU(2) gauge fields*

$$\sin \theta_W = \frac{e}{g}, \quad \sin^2 \theta_W = 0.2230 \pm 0.0004$$

- *The Fermi interaction*



$$G_F = \frac{\sqrt{2}g^2}{8M_W^2} = 1.16639 \pm 0.00001 \times 10^{-5} \text{ GeV}^{-2}$$

Neutral weak current

- *Electromagnetic current*

$$J_{EM}^\mu = \frac{2}{3}\bar{u}\gamma^\mu u - \frac{1}{3}\bar{d}\gamma^\mu d - \frac{1}{3}\bar{s}\gamma^\mu s$$

- *Neutral current*

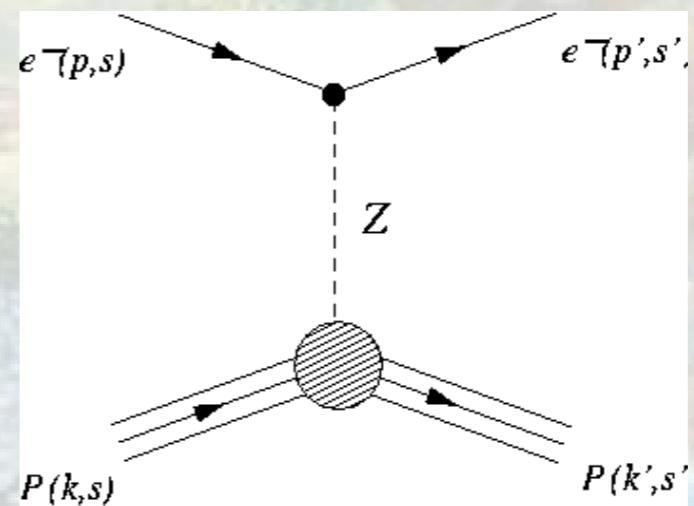
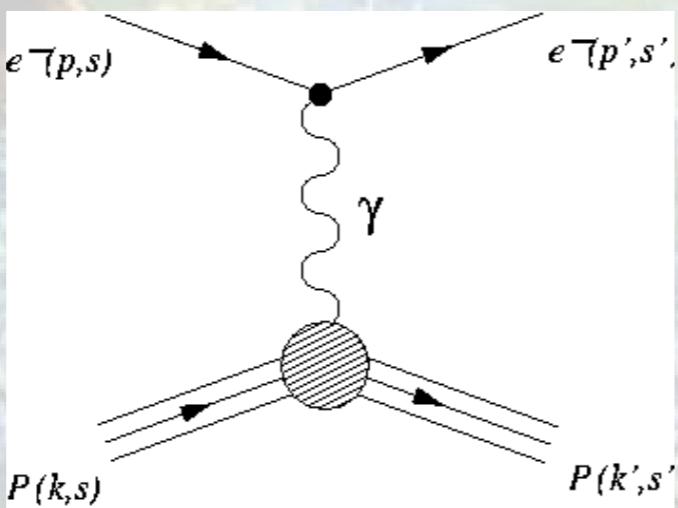
$$\begin{aligned} J_{NC}^\mu = & \bar{u}\gamma^\mu \left(\frac{1}{2} - \frac{4}{3}\sin^2\theta_W - \frac{1}{2}\gamma_5 \right) u + \bar{d}\gamma^\mu \left(-\frac{1}{2} + \frac{2}{3}\sin^2\theta_W + \frac{1}{2}\gamma_5 \right) d \\ & + \bar{s}\gamma^\mu \left(-\frac{1}{2} + \frac{2}{3}\sin^2\theta_W + \frac{1}{2}\gamma_5 \right) s \end{aligned}$$

- *Charged current*

- *More complicated: CKM mixing matrix*

Experimental Strategy

- Measure the interference between photon and neutral currents in parity-violating electron scattering (PVES)



$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim \frac{\left| \begin{array}{c} \nearrow \gamma \\ \searrow \gamma \end{array} \right\rangle \left\langle \begin{array}{c} Z^0 \\ \searrow \gamma \end{array} \right|}{\left| \begin{array}{c} \nearrow \gamma \\ \searrow \gamma \end{array} \right\rangle \left\langle \begin{array}{c} \nearrow \gamma \\ \searrow \gamma \end{array} \right|} \sim \frac{|M_{PV}^{NC}|}{|M^{EM}|} \sim \frac{Q^2}{(M_Z)^2}$$

PVES from a proton target

- *Total asymmetry*

$$A^{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} = \left[\frac{-G_F Q^2}{\pi \alpha \sqrt{2}} \right] \frac{\varepsilon G_E^{p\gamma} G_E^{pZ} + \tau G_M^{p\gamma} G_M^{pZ} - \frac{1}{2}(1 - 4 \sin^2 \theta_W) \varepsilon' G_M^{p\gamma} \tilde{G}_A^p}{\varepsilon (G_E^{p\gamma})^2 + \tau (G_M^{p\gamma})^2}$$

Charge symmetry: neglect u-d mass difference, EM corrections invariance under u-d exchange

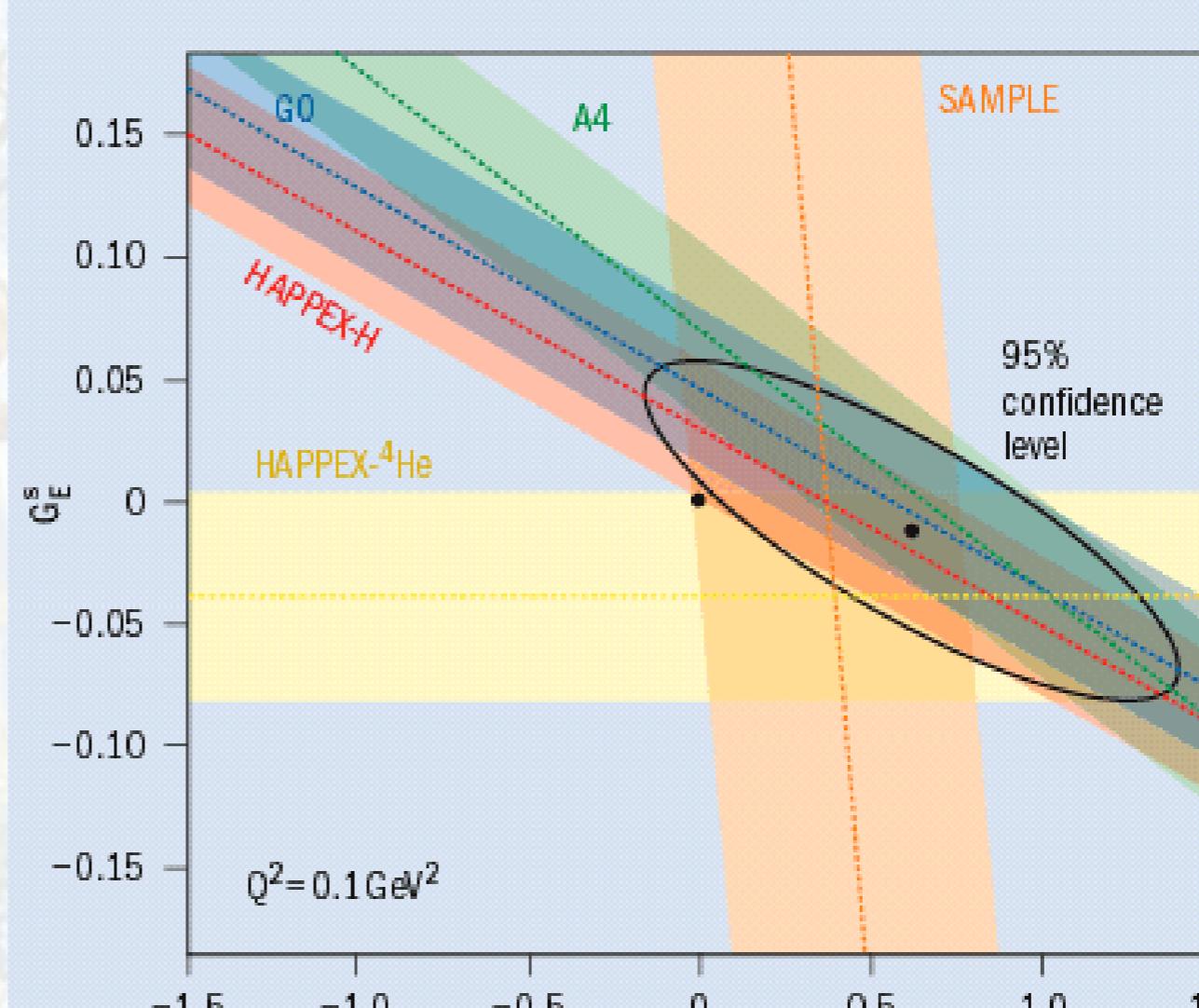
$$G_{E,M}^{p(u)} = G_{E,M}^{n(d)}, \quad G_{E,M}^{p(s)} = G_{E,M}^{n(s)}, \dots$$

$$4G_{E,M}^{pZ} = (1 - 4 \sin^2 \theta_W) G_{E,M}^{p\gamma} - G_{E,M}^{n\gamma} - G_{E,M}^s$$

Strangeness

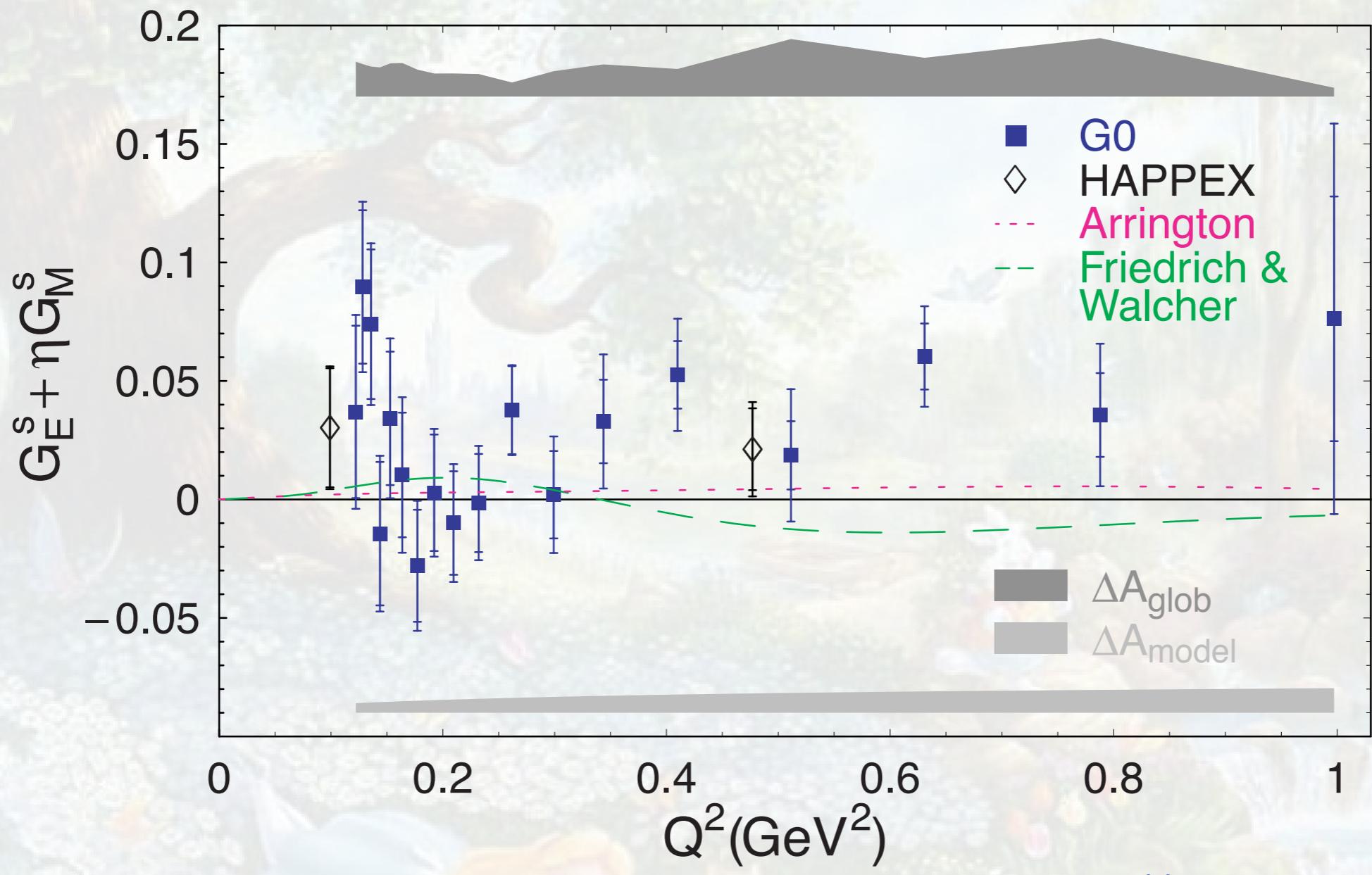
Experimental Programs

- **SAMPLE Collaboration at MIT-Bates**
 - Backward angle scattering: H & D @ $Q^2 \sim 0.1 \text{ GeV}^2$
- **HAPPEX Experiment at Jefferson Lab**
 - Forward angle: H @ $Q^2 \sim 0.5, 0.1 \text{ GeV}^2$; He4 @ $Q^2 \sim 0.1 \text{ GeV}^2$
- **PVA4 at Mainz**
 - Forward angle: H @ $Q^2 \sim 0.23, 0.1 \text{ GeV}^2$
 - Backward run: H @ $Q^2 \sim 0.23 \text{ GeV}^2$ (*in progress*)
- **GO Experiment at Jefferson Lab**
 - Forward angle: H @ $Q^2 \sim 0.12 - 1.0 \text{ GeV}^2$
 - Backward run: H & D @ $Q^2 \sim 0.23, 0.6 \text{ GeV}^2$ (*in progress*)

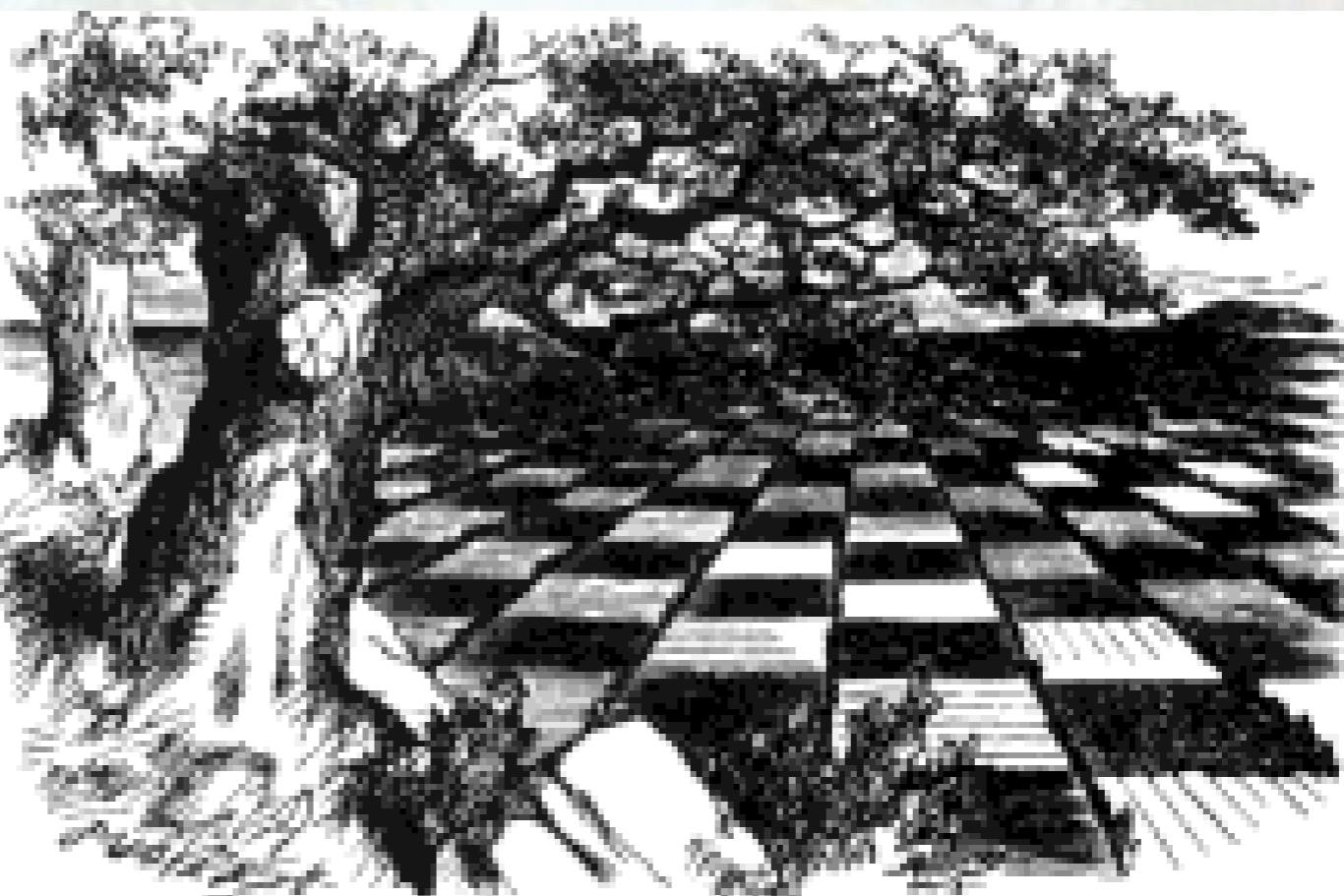


CERN Courier

*Suggestion of a large & positive
strangeness magnetic moment*



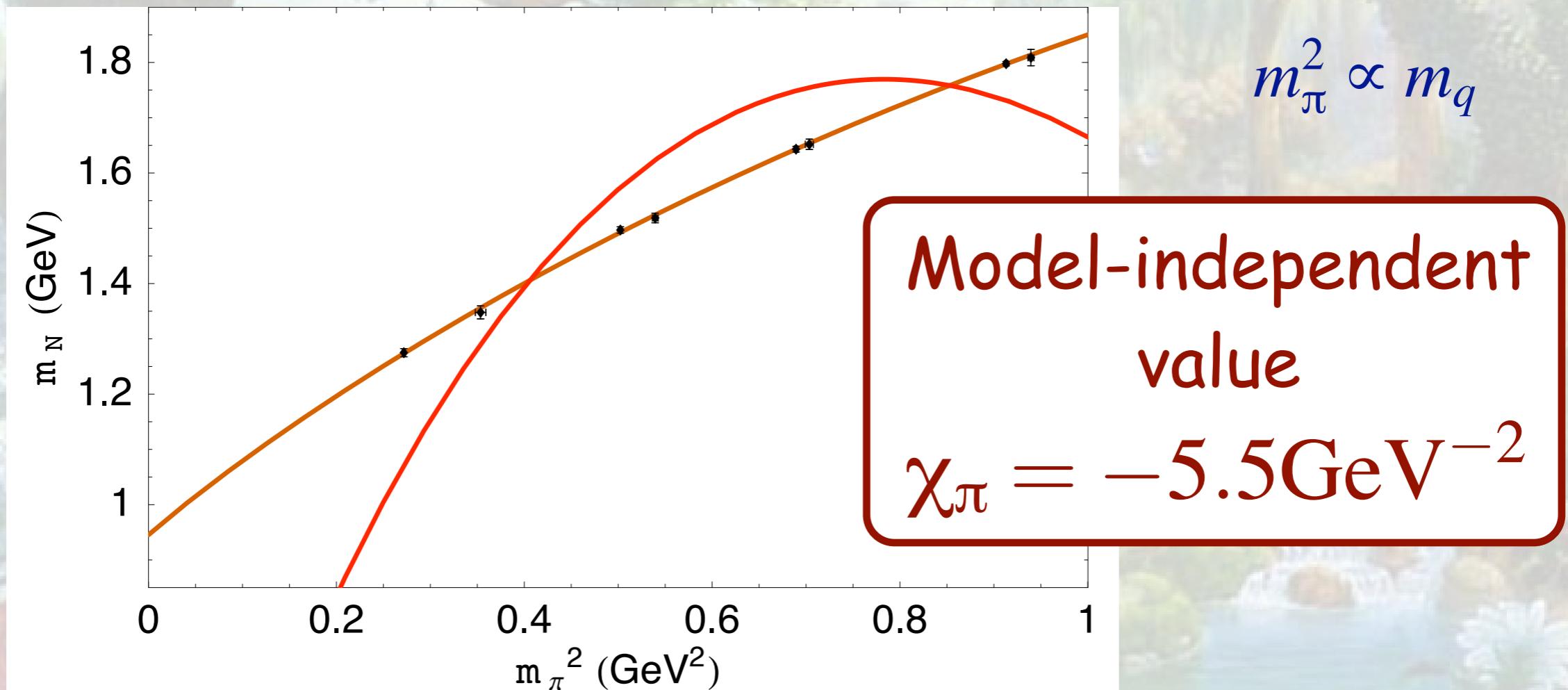
GO Collaboration



THEORY:
Strangeness content in a QCD-based approach.

Chiral Extrapolation

ChPT: $M_N = c_0 + c_2 m_\pi^2 + \chi_\pi m_\pi^3 + \dots$



$$\chi_\pi \simeq -0.63 \text{ GeV}^{-2}$$

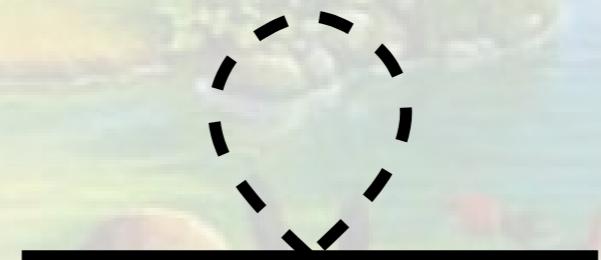
Finite-range regularisation


$$= -\frac{3g_A^2}{16\pi^2 f_\pi^2} \int_0^\infty dk \frac{k^4}{k^2 + m_\pi^2}$$

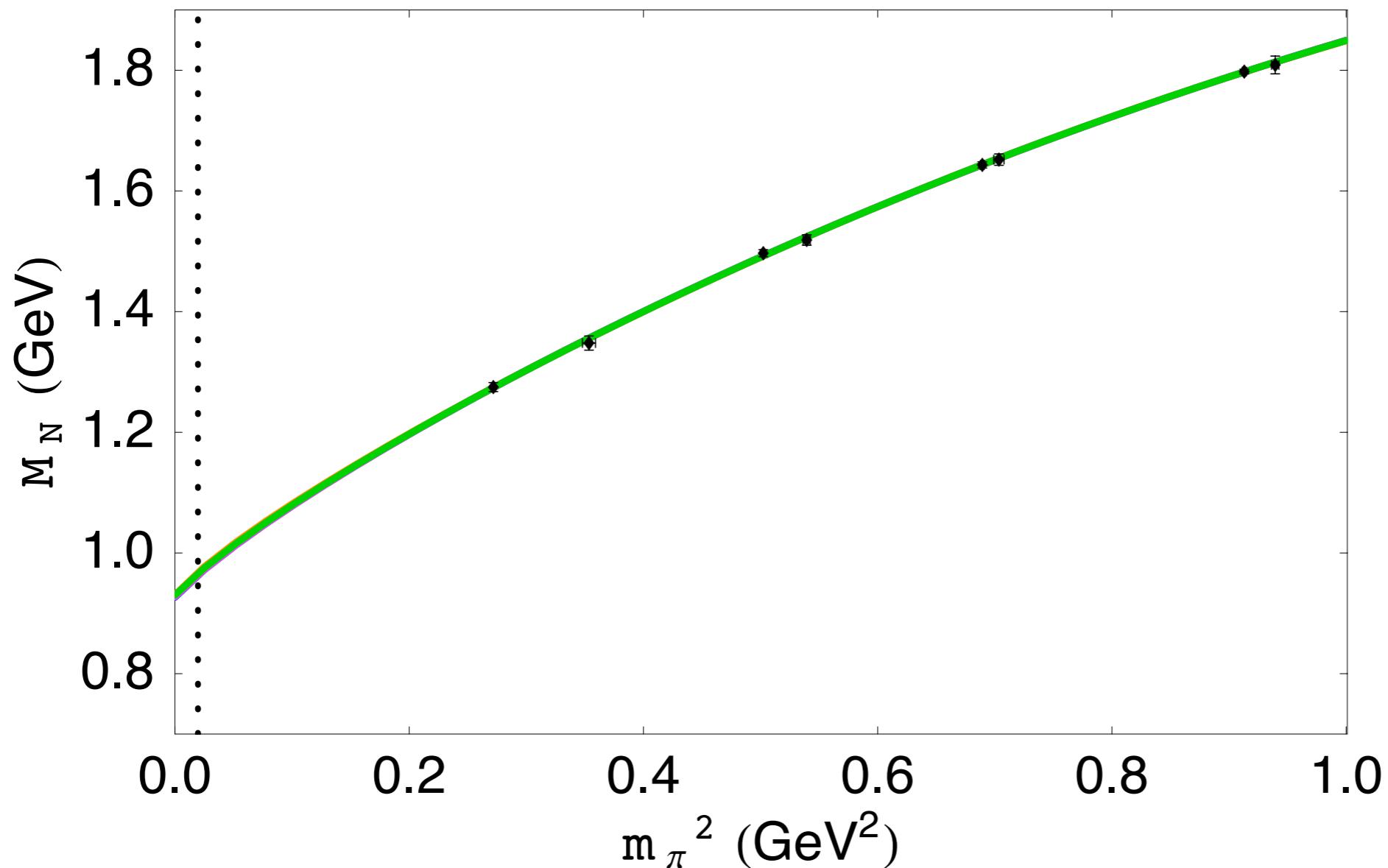
Use functional form to regularise ultra-violet behaviour of loop integral, eg. dipole

$$-\frac{3g_A^2}{16\pi^2 f_\pi^2} \int_0^\infty dk \frac{k^4}{k^2 + m_\pi^2} \left(\frac{\Lambda^2}{\Lambda^2 + k^2} \right)^4$$

Exactly the same model-independent nonanalytic contribution!



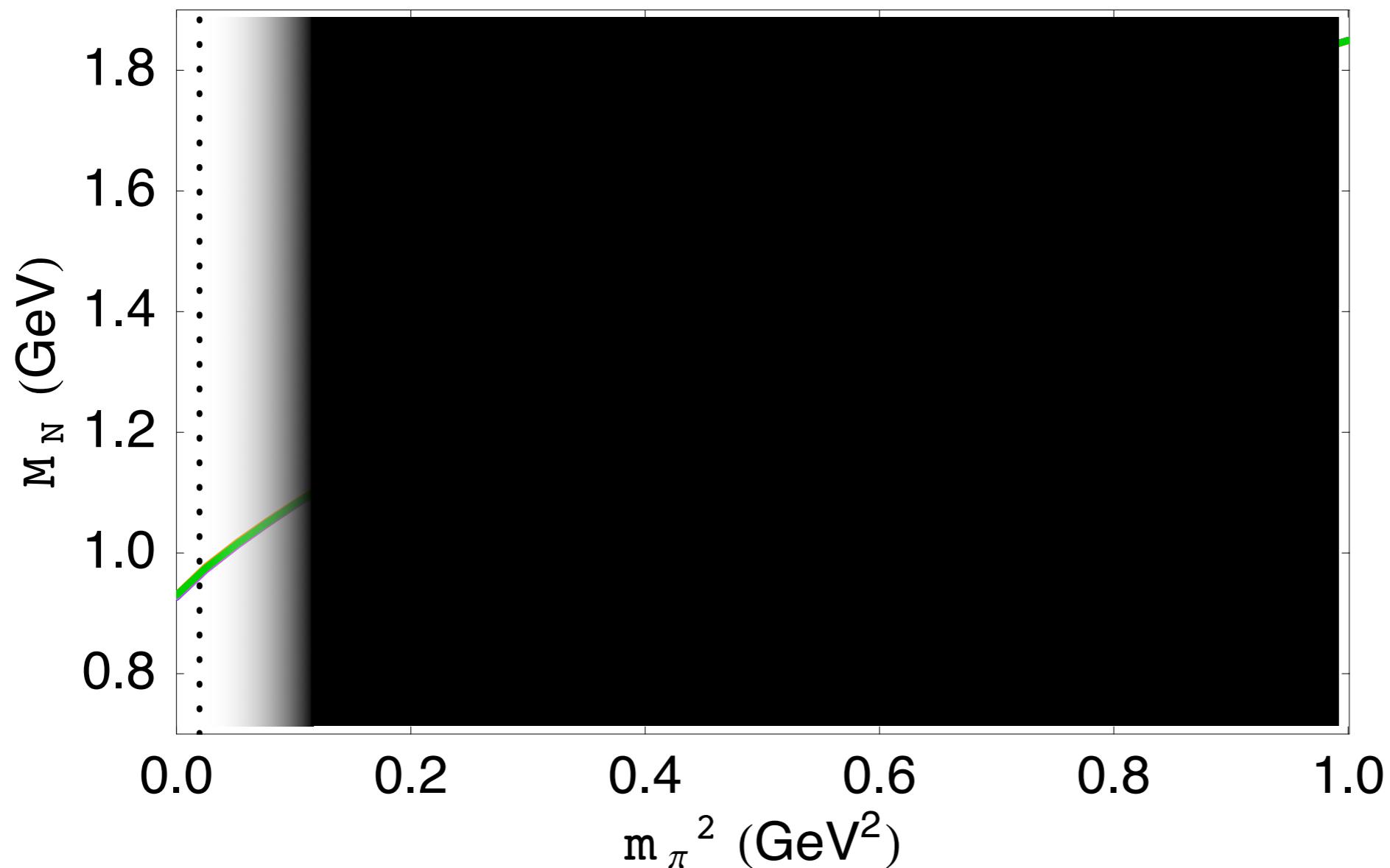
Finite-range regularisation.. 2



Dipole
Theta
Monopole
Gaussian

Independent-of-model chiral extrapolation

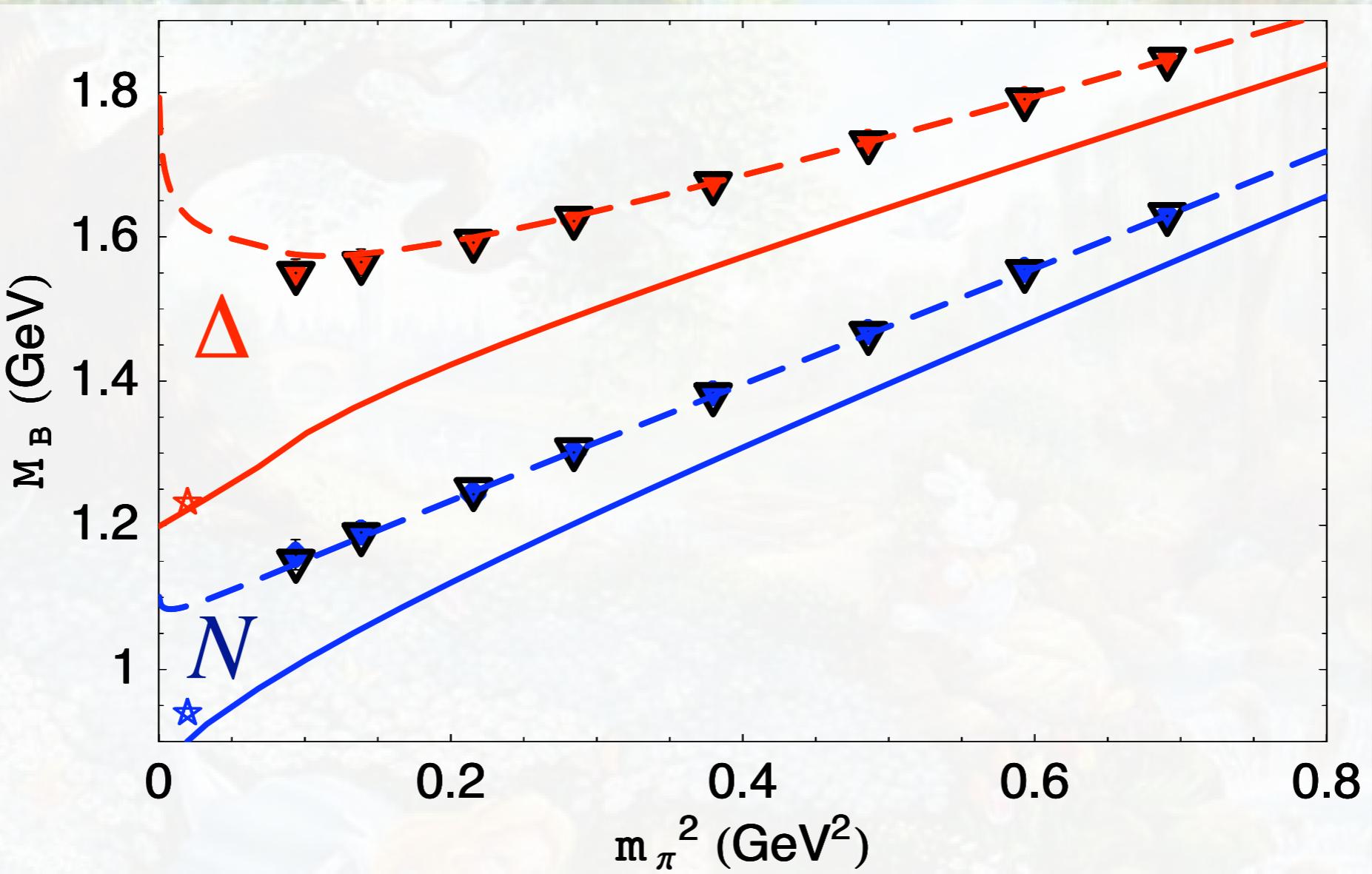
Finite-range regularisation.. 2



Dipole
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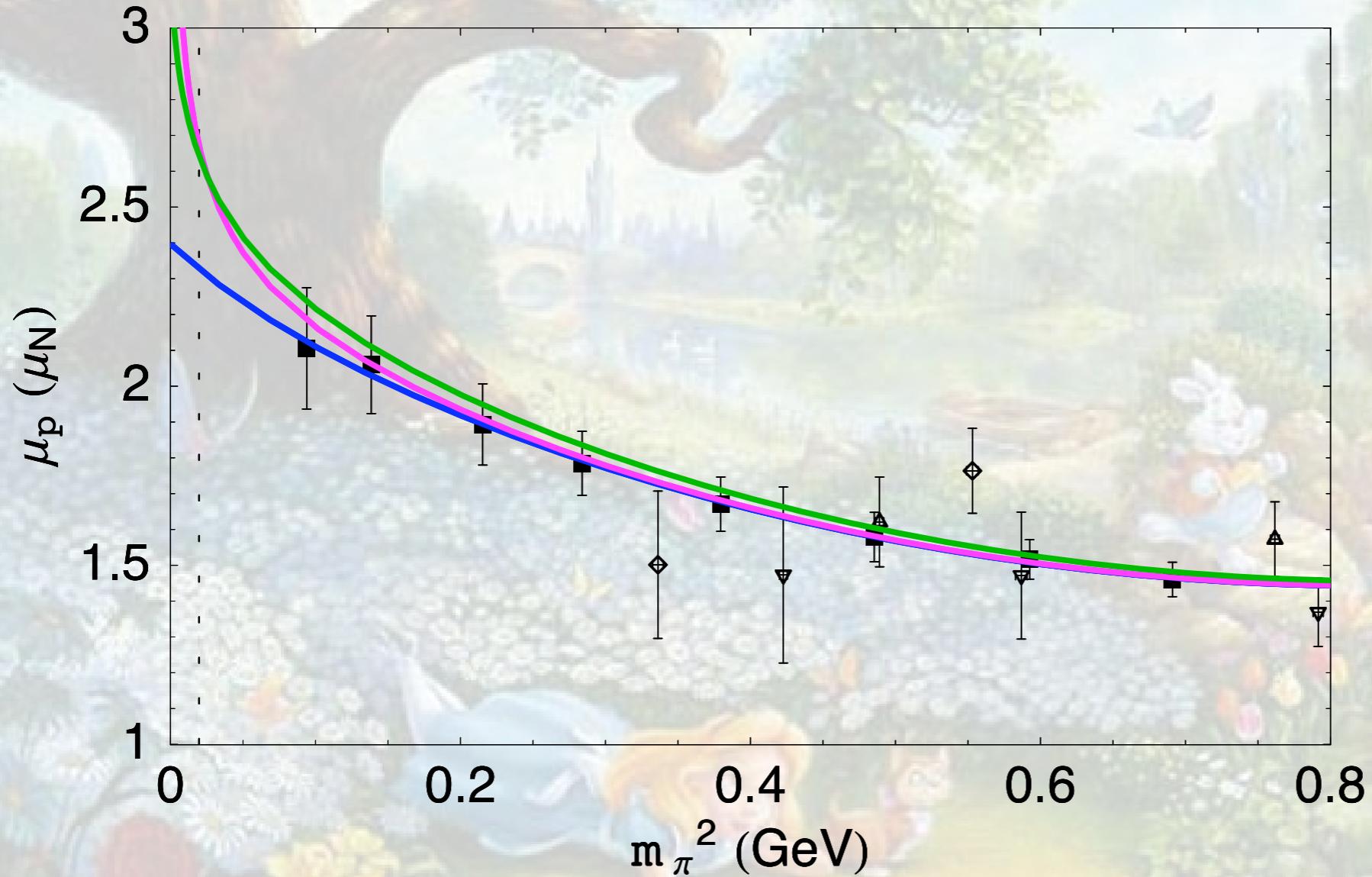
Baryon Masses



*Establishing a connection between quenched
simulations and QCD*

Differences described by pion-loop corrections

Proton Magnetic Moment



Finite-volume
Quenched
QCD

Apply unquenching technique to EM structure

Strangeness

Assume charge symmetry

$$p = \frac{2}{3}u^p - \frac{1}{3}u^n + O_N$$

$$n = -\frac{1}{3}u^p + \frac{2}{3}u^n + O_N$$

$$3O_N = 2p + n - u^p$$

$$\Sigma^+ = \frac{2}{3}u^\Sigma - \frac{1}{3}s^\Sigma + O_\Sigma$$

$$\Sigma^- = -\frac{1}{3}u^\Sigma - \frac{1}{3}s^\Sigma + O_\Sigma$$

$$\Sigma^+ - \Sigma^- = u^\Sigma$$

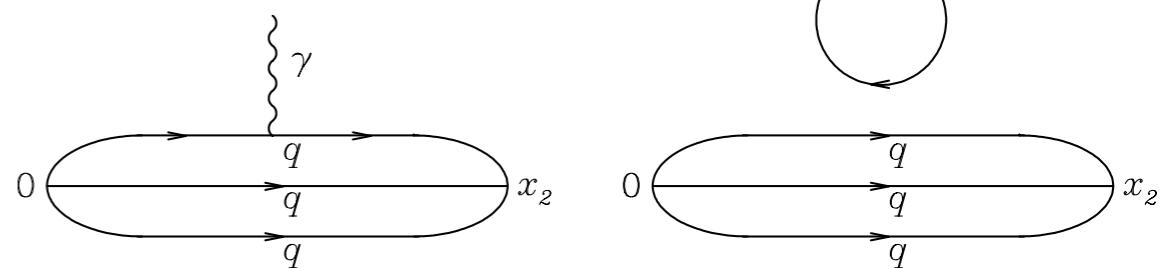
$$3O_N = 2p + n - \frac{u^p}{u^\Sigma}(\Sigma^+ - \Sigma^-)$$

Lattice QCD

$$3O_N = p + 2n - u^n$$

$$\Xi^0 - \Xi^- = u^\Xi$$

$$3O_N = p + 2n - \frac{u^n}{u^\Xi}(\Xi^0 - \Xi^-)$$



The loop part

$$O_N = \text{---} \overset{\times}{\textcircled{*}} u, d, s \text{---}$$

“ l ” loop contribution

$$= \frac{2}{3} {}^l G_M^u - \frac{1}{3} {}^l G_M^d - \frac{1}{3} {}^l G_M^s$$

$$\begin{aligned} O_N &= -\frac{1}{3} ({}^l G_M^d + {}^l G_M^s) \\ &= \frac{{}^l G_M^s}{3} \left(1 - \frac{{}^l R_d^s}{{}^l R_d^u} \right) \end{aligned}$$

$${}^l G_M^u = {}^l G_M^d$$

QCD equality for $m_u = m_d$

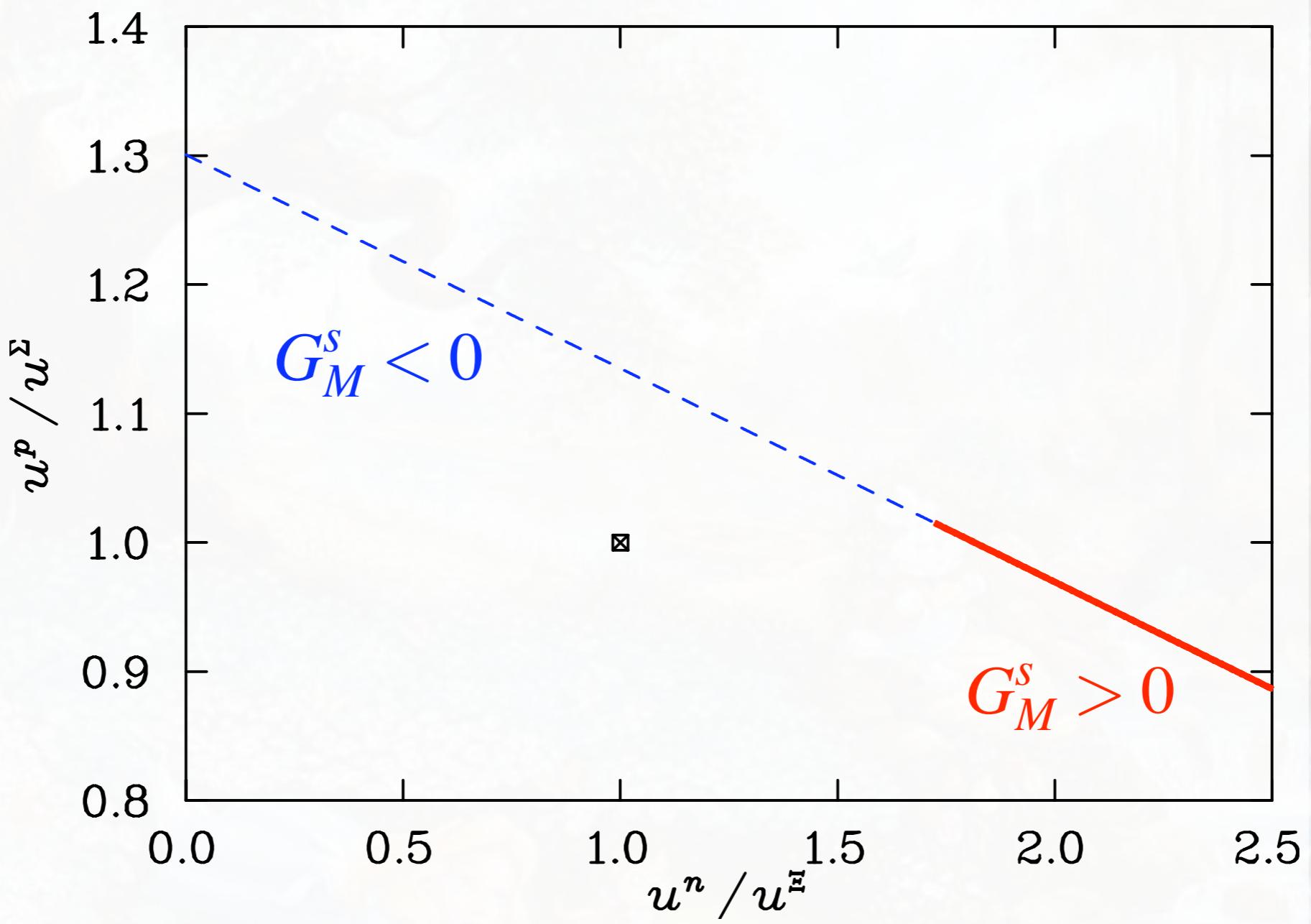
$${}^l R_d^s = {}^l G_M^s / {}^l G_M^d = 0.139 \pm 0.042$$

Kaon strength relative to pion

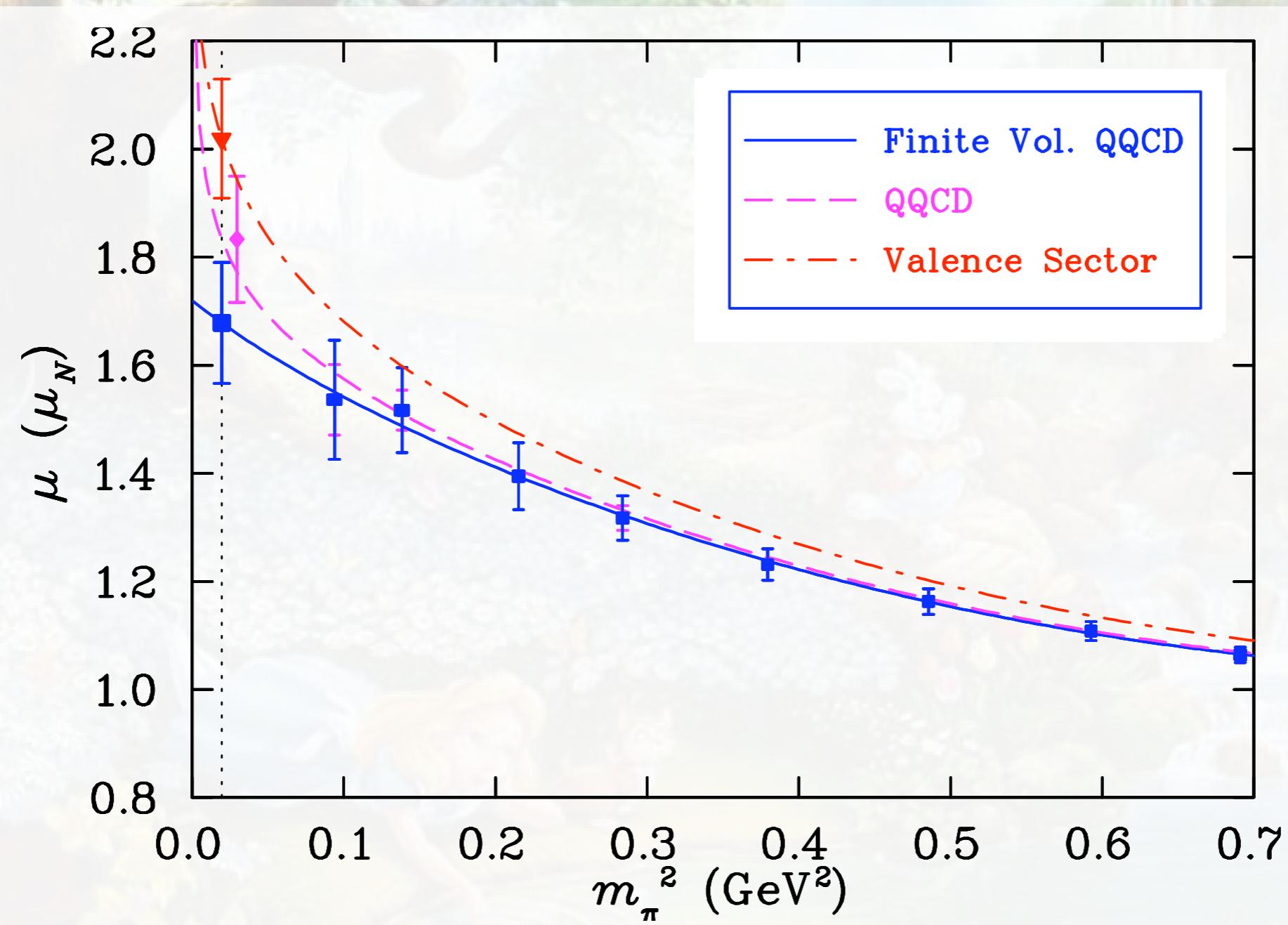
Strangeness Magnetic Moment

$$G_M^s = \left(\frac{{}^l R_d^s}{1 - {}^l R_d^s} \right) \left[2p + n - \frac{u^p}{u^\Sigma} (\Sigma^+ - \Sigma^-) \right]$$
$$G_M^s = \left(\frac{{}^l R_d^s}{1 - {}^l R_d^s} \right) \left[p + 2n - \frac{u^n}{u^\Xi} (\Xi^0 - \Xi^-) \right]$$

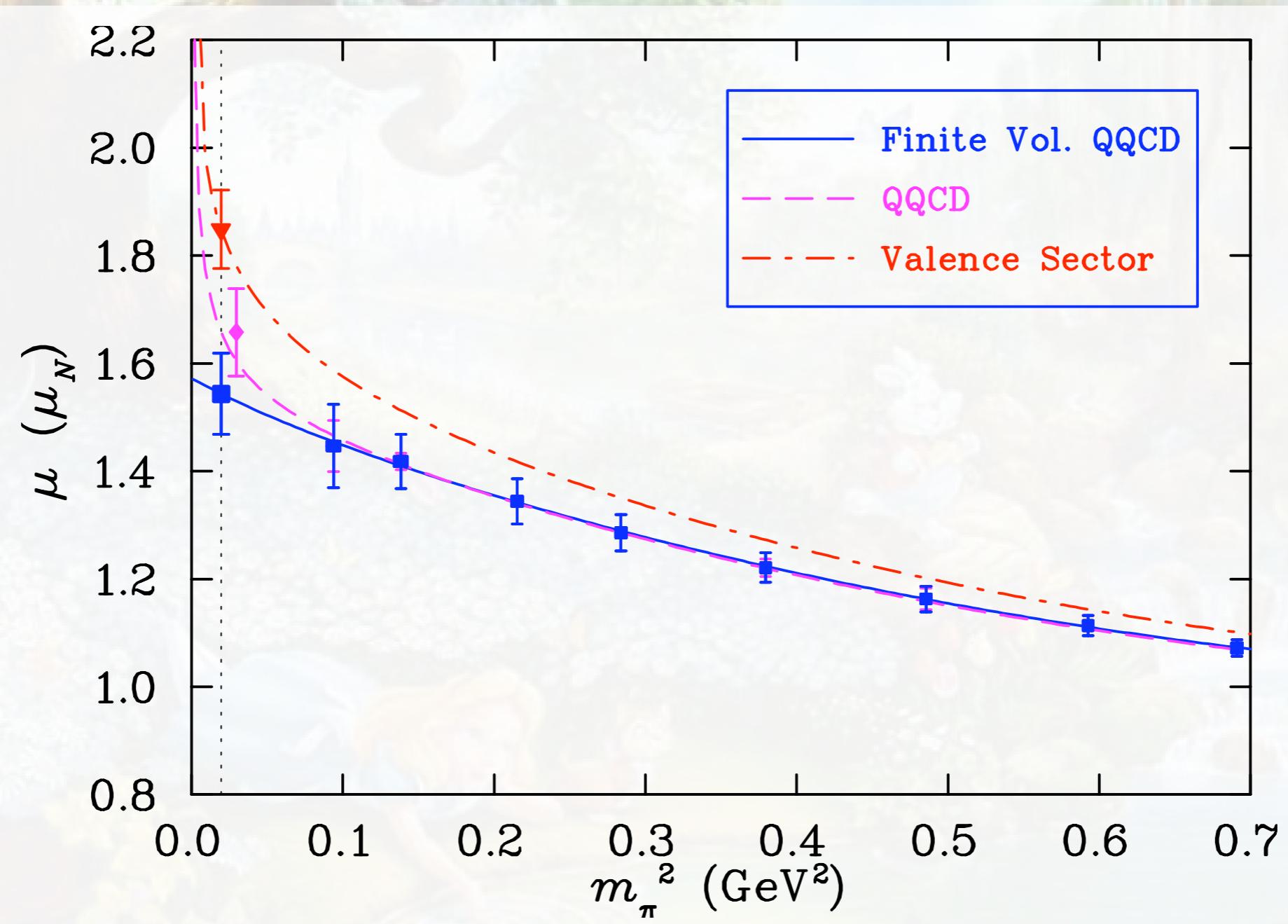
*Two equations give linear constraint between
p/Sigma and n/Xi ratios*



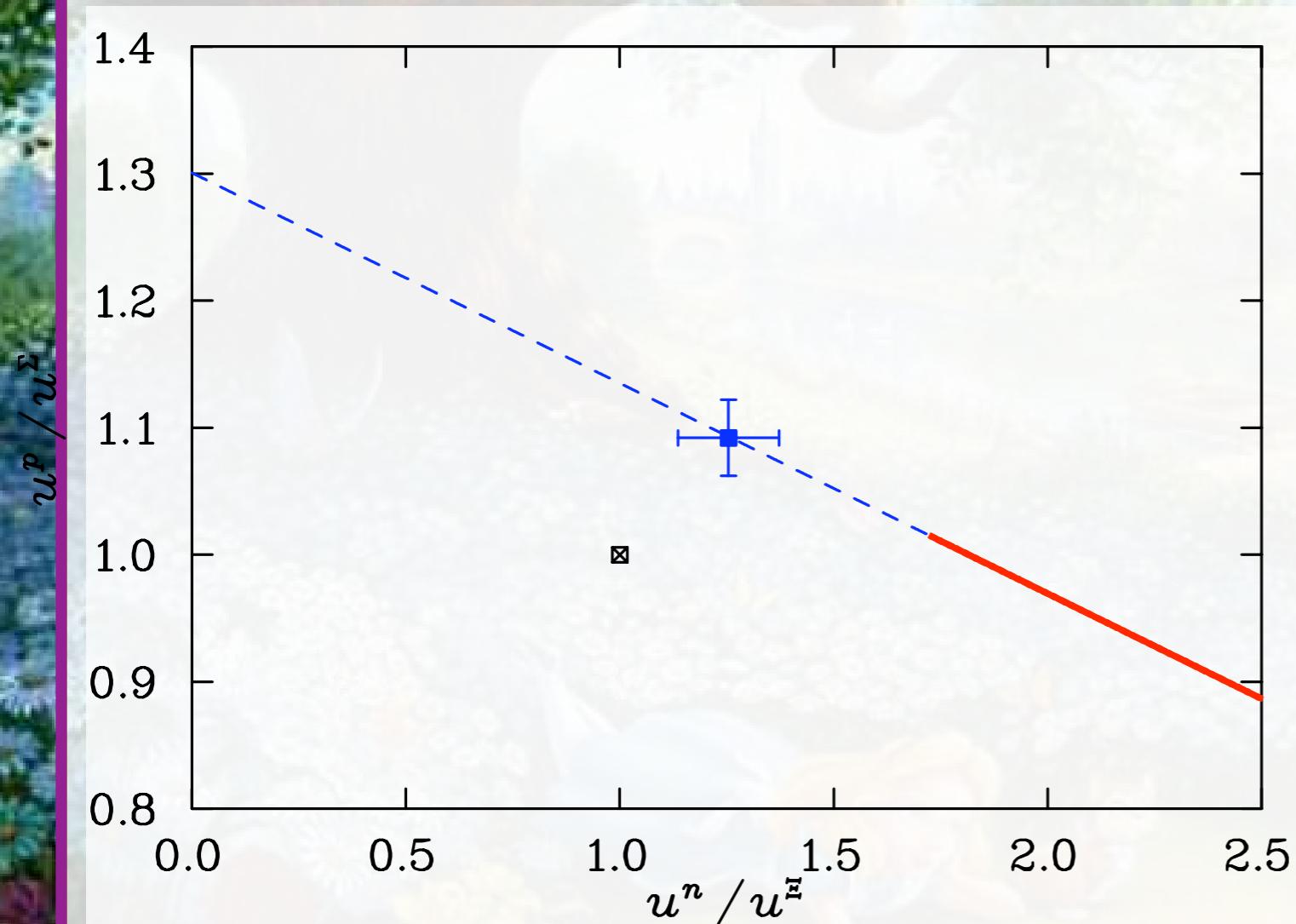
u-quark in the proton



u-quark in the Sigma



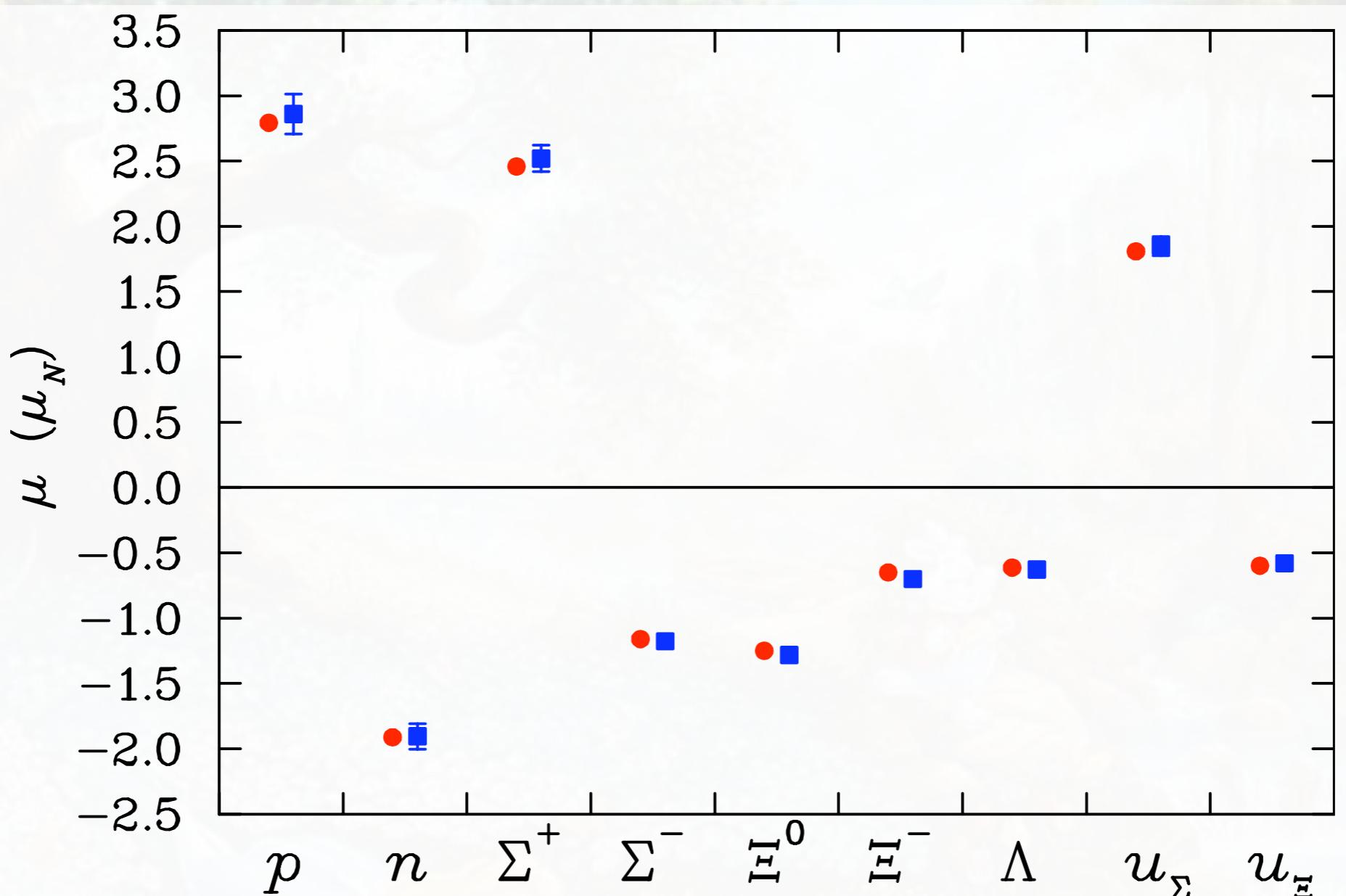
Final Result



$$\frac{u^p}{u^\Sigma} = 1.092 \pm 0.030$$

$$\frac{u^n}{u^\Xi} = 1.254 \pm 0.124$$

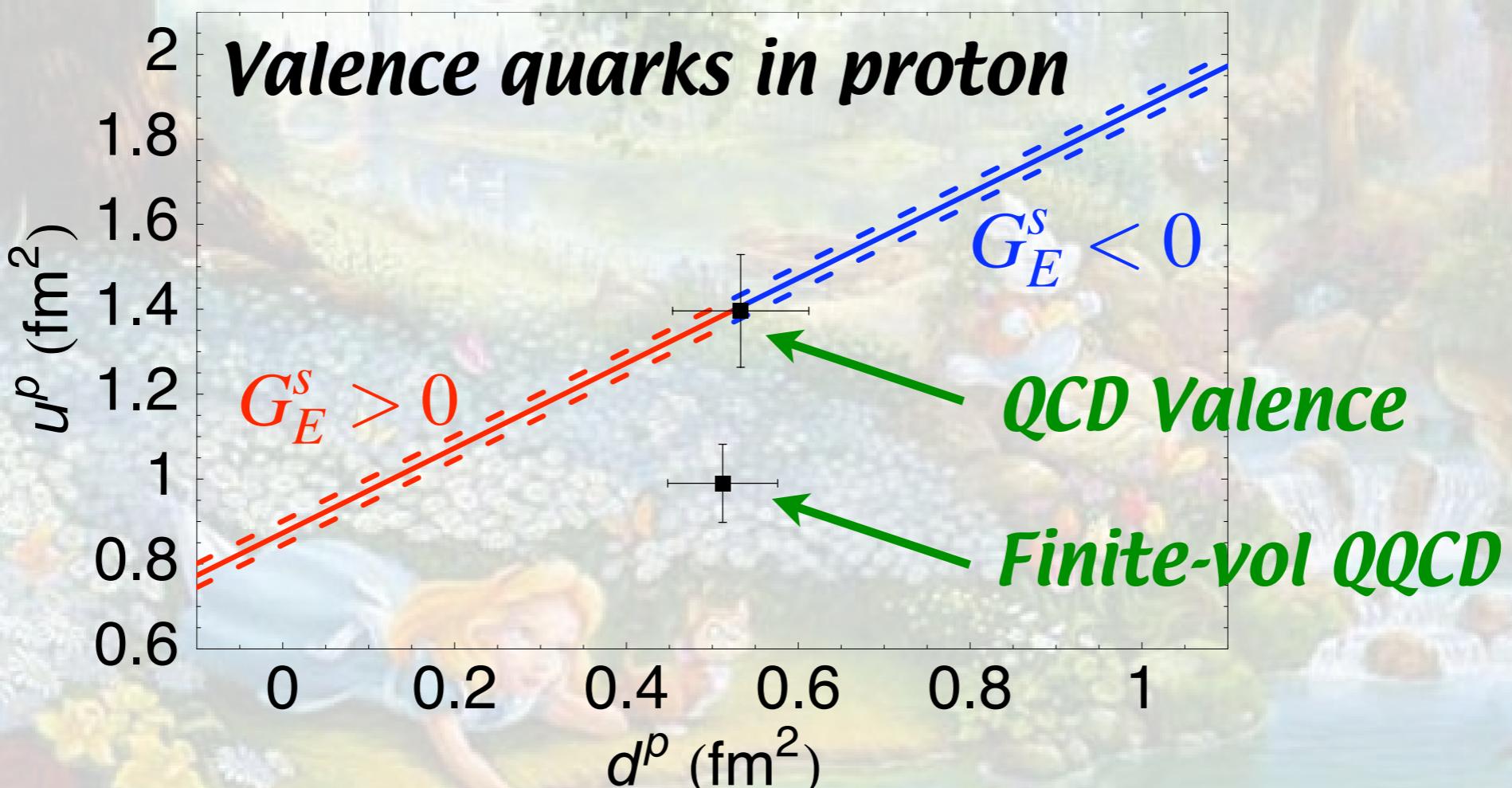
$$G_M^s = -0.046 \pm 0.022 \mu_N$$



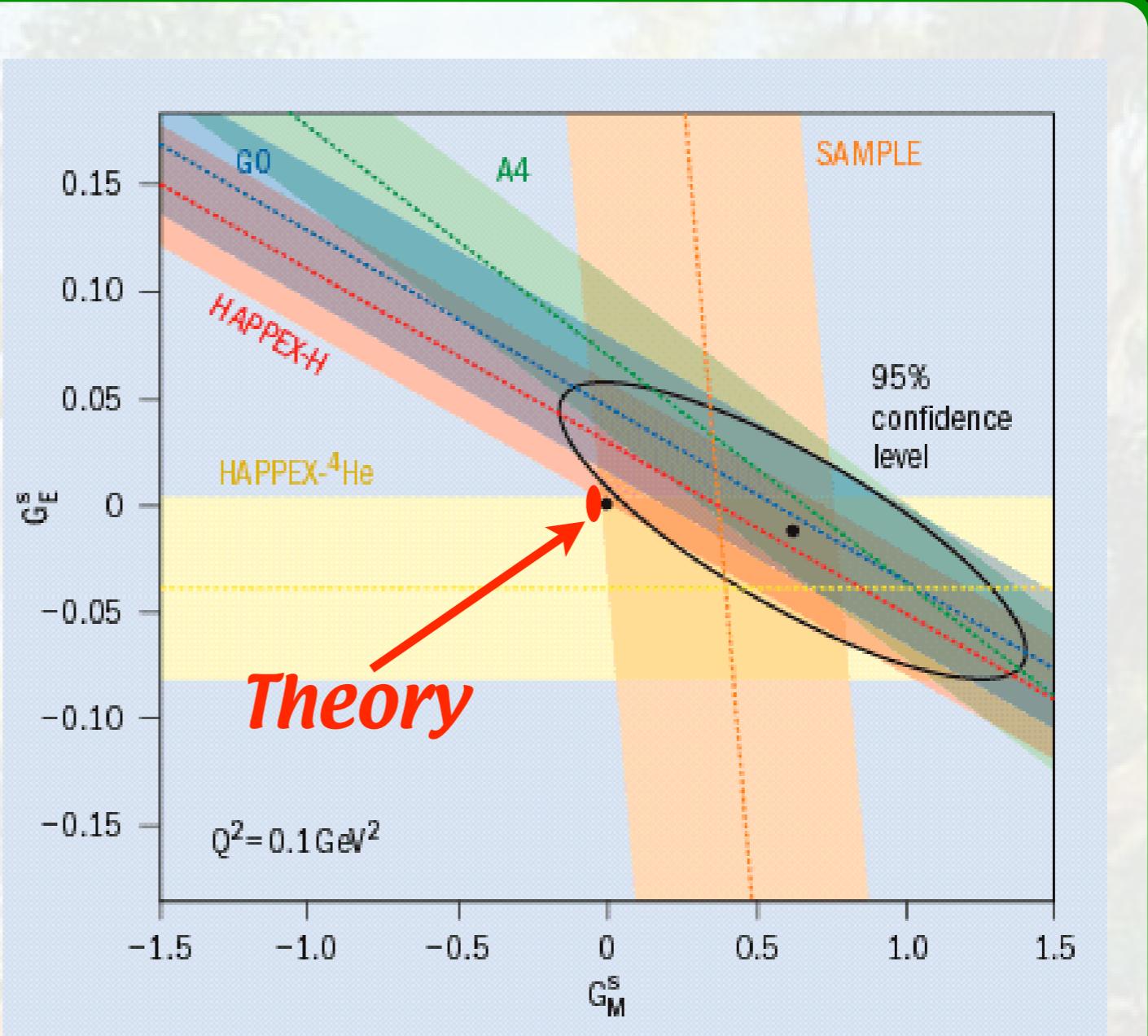
Hyperon magnetic moments

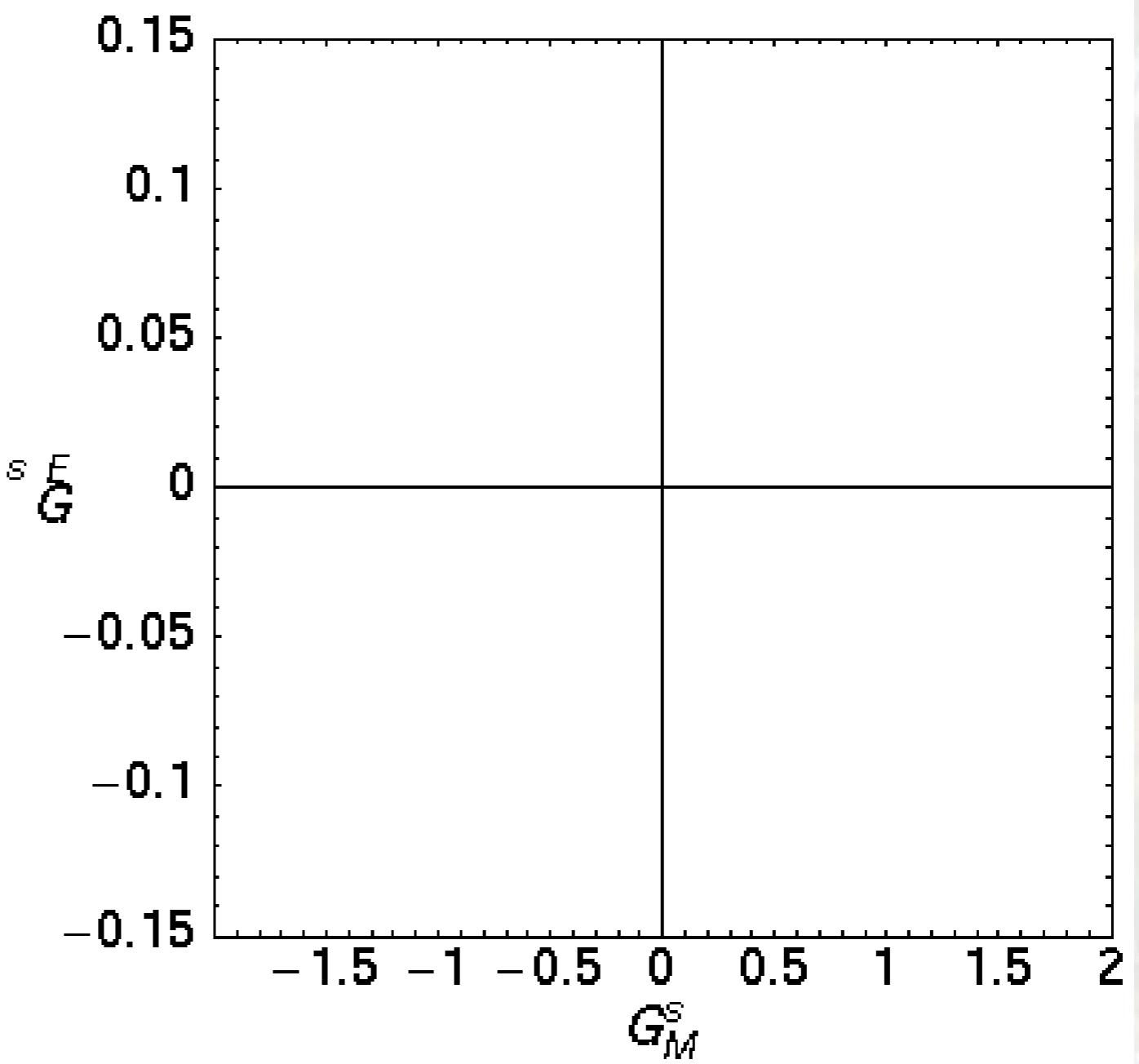
Electric Form Factor

Limit knowledge of hyperon charge radii

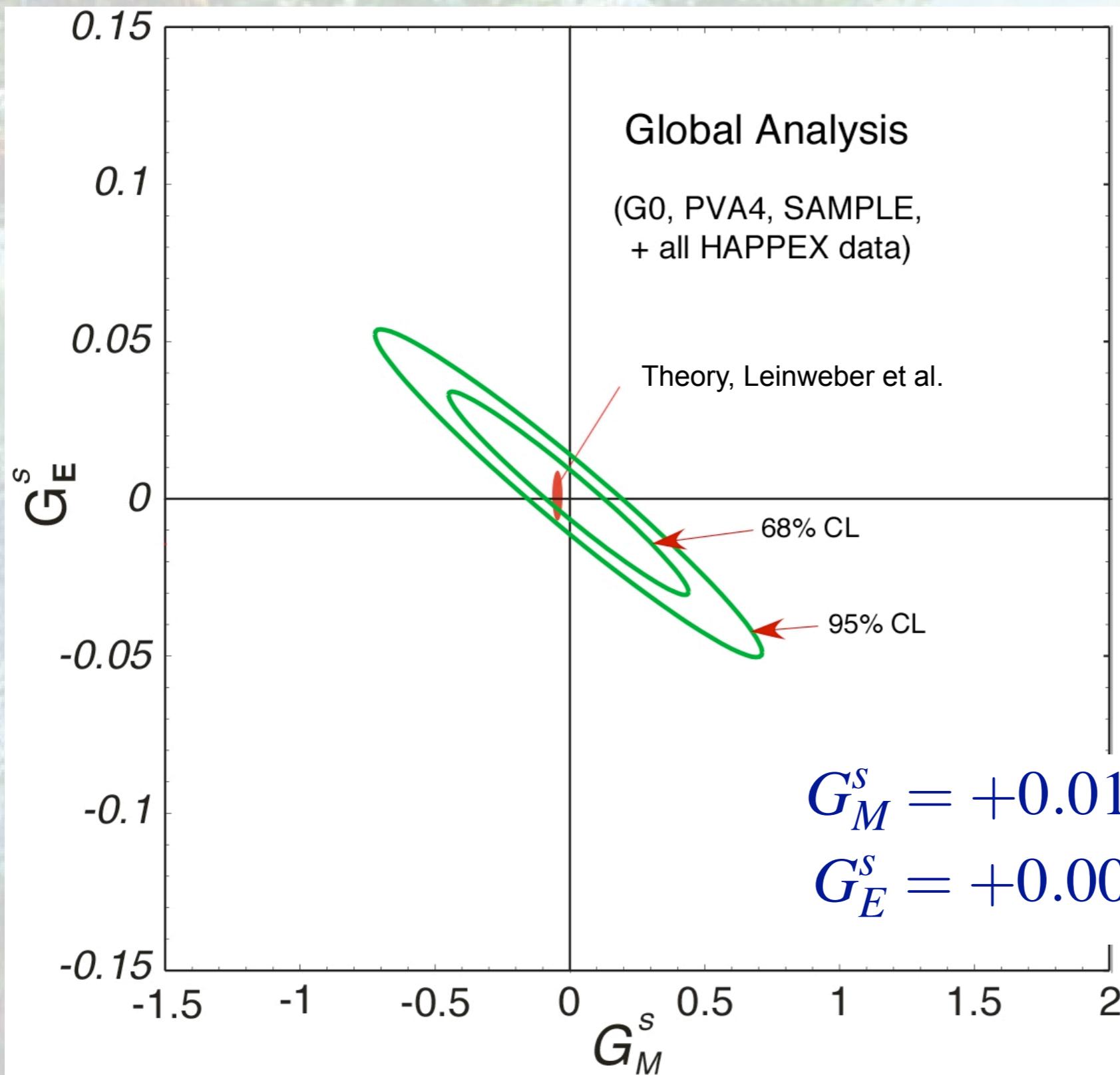


$$G_E^s(Q^2 = 0.1) = +0.001 \pm 0.004 \pm 0.004$$





*New 2006 HAPPEX
results + global
analysis of all
world data*



Electron scattering can probe charge and magnetisation distributions in the nucleon

Chiral symmetry important in low-energy QCD

Chiral symmetry provides vital information on nucleon structure in QCD

Weak-interaction violates parity symmetry, allows flavour separation of nucleon structure

Theoretical approach providing a significant advance in the understanding of the nucleon



THANKS!